

ON JOINT OBSERVABLE FOR F-QUANTUM SPACES

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1. F-quantum spaces

There are many analogies between fuzzy set theory and quantum logic models of quantum mechanics as it has been noticed in [7]. In a new axiomatic model for measurement of quantum mechanical observables based on fuzzy sets ideas we solve the problem of existence of a joint observable for a given system of observables of an F-quantum space.

Definition 1.1. By an F-quantum space we mean a couple (X, M) , where X is a non-empty set and M is a subset of $[0, 1]^X$ with (i) if $1(x) = 1$ for any $x \in X$, then $1 \in M$; (ii) if $f \in M$, then $1-f \in M$; (iii) if $1/2(x) = 1/2$ for any $x \in X$, then $1/2 \in M$; (iv) if $\{f_n\} \subset M$, then $\bigvee_n f_n := \sup_n f_n \in M$. Using

Using the terminology of Piasecki [4] M is a soft fuzzy \mathcal{G} -algebra of fuzzy sets of X . Moreover, $\bigwedge_n f_n := \inf_n f_n \in M$ for any $\{f_n\} \subset M$. A non-empty subset $\mathcal{M} \subset M$ is said to be a Boolean algebra (\mathcal{G} -algebra) of an F-quantum space (X, M) if (i) there are the minimal and maximal elements $0_{\mathcal{M}}, 1_{\mathcal{M}} \in \mathcal{M}$ such that, for any $f \in \mathcal{M}$, $0_{\mathcal{M}} \leq f \leq 1_{\mathcal{M}}$; (ii) a mapping $\perp : f \mapsto f^\perp = 1-f$, $f \in \mathcal{M}$, satisfies $f \vee f^\perp = 1_{\mathcal{M}}$ for any $f \in \mathcal{M}$; (iii) \mathcal{M} is with respect to $\wedge, \vee, \perp, 0_{\mathcal{M}}, 1_{\mathcal{M}}$ a Boolean algebra (\mathcal{G} -algebra). We note that $0_{\mathcal{M}} \neq 1_{\mathcal{M}}$, in the opposite case $0_{\mathcal{M}} = 1/2 \in M$. It is simple that M is a Boolean \mathcal{G} -algebra iff $f = f^2$ for any $f \in M$, i.e., M consists from crisp sets.

Let (Ω, \mathcal{A}) be a measurable space. We say that a mapping x :

$\mathcal{A} \rightarrow M$ is an \mathcal{A} -observable (\mathcal{A} - σ -observable) of (X, M) if (i) $x(A) + x(A') = 1$ for any $A \in \mathcal{A}$ / here A' denotes the complement of A in Ω /; (ii) $x(\bigcup_i A_i) = \bigvee_i x(A_i)$ for any finite (countable) number of sets from \mathcal{A} . For the quantum mechanics it is of a great importance the case when \mathcal{A} is the Borel σ -algebra of some separable Banach space Y , in particular, when $Y = R_1$. In this case we call x an observable. The range of an (σ -) observable x , that is, the set $\mathcal{Q}(x) = \{x(A) : A \in B(R_1)\}$ is a Boolean algebra (σ -algebra) of (X, M) with the minimal and maximal elements $x(\emptyset)$ and $x(R_1)$, respectively.

We say that a system $\{x_t : t \in T\}$, where x_t is a (σ -) observable of (X, M) , has a joint (σ -) observable if there is an $R(T)$ -observable ($B(T)$ - σ -observable) x of (X, M) , where $R(T) (B(T))$ is the minimal algebra (σ -algebra) of R_1^T containing all finite-dimensional rectangles, such that $x(\bigcup_{t \in \alpha} A_t) = \bigvee_{t \in \alpha} x_t(A_t)$ for any $A_t \in B(R_1)$ and any finite non-empty $\alpha \subset T$.

2. Joint observables

Theorem 2.1. Let $\{x_t : t \in T\}$ be a system of (σ -) observables of an F-quantum space (X, M) . The following assertions are equivalent: (i) $x_s(\emptyset) = x_t(\emptyset)$ for any $s, t \in T$.

(ii) $x_s(R_1) = x_t(R_1)$ for any $s, t \in T$.

(iii) Any subsystem $\{x_s, x_t\}$, $s, t \in T$ has a joint (σ -) observable.

(iv) $\{x_t : t \in T\}$ has a joint (σ -) observable.

Proof. The implication (iii) \rightarrow (i) is evident. Conversely; for simplicity we put $x = x_s$, $y = x_t$. Denote by $\mathcal{D} = \{A \times B : A, B \in B(R_1)\}$ and define a mapping $h : \mathcal{D} \rightarrow M$ via $h(A \times B) = x(A) \wedge y(B)$, $A, B \in B(R_1)$. Using the standard techniques we may show that h is a \mathcal{D} -observable of (X, M) . Analogically may be proved the implication (iii) \rightarrow (iv).

For the case of σ -observables, the proof of (i) \rightarrow (iv).

Since any Boolean algebra is isomorphic to some algebra of subsets [8], the conditions (i) - (iv) for observables are equivalent in order to exist a Boolean algebra $\mathcal{M} \subset M$ including all ranges $\mathcal{R}(x_t)$. Using the Zorn lemma, we may show that there is a maximal Boolean algebra \mathcal{M}_0 of (X, M) containing \mathcal{M} . We assert that \mathcal{M}_0 is a Boolean σ -algebra of (X, M) . Indeed, let $\{f_n\} \subset \mathcal{M}_0$. Then $f = \bigvee_{n=1}^{\infty} f_n$ is an element of M , and $f \vee (1-f) = x_t(R_1)$ for any $t \in T$. In fact, $f \vee (1-f) = \max\{f, 1-f\} = (1 + |2f-1|)/2 = \lim_n (1 + |2g_n-1|)/2 = x_t(R_1)$, where $g_n = \bigvee_{i=1}^n f_i \in \mathcal{M}_0$, $n \geq 1$, and $g_n \vee (1-g_n) = x_t(R_1)$, $n \geq 1$. Define a Boolean algebra $\mathcal{M}_f = \{x_t(\emptyset), x_t(R_1), f, 1-f\}$ of (X, M) . The maximality of \mathcal{M}_0 entails $\mathcal{M}_f \subset \mathcal{M}_0$.

The σ -distributivity of \mathcal{M}_0 entails that \mathcal{M}_0 has a σ -strong extension property, for definition see [8]. According to Sikorski [8, Theorem 37.1], the joint observable x of X, M generated by $\{x_t: t \in T\}$ may be extended uniquely to a joint- σ -observable of (X, M) . Q.E.D.

Let $f: R_1 \rightarrow R_1$ be a Borel measurable function and let x be a σ -observable of (X, M) . Then by $f \circ x$ we mean a σ -observable of (X, M) defined by $f \circ x(E) = x(f^{-1}(E))$, $E \in B(R_1)$. According to the terminology of the theory of quantum logics (see [9]), we say that two σ -observables of (X, M) x and y are compatible if $x(\emptyset) = y(\emptyset)$.

Theorem 2.2. Let $\{x_t: t \in T\}$ be a system of σ -observables of an F-quantum space (X, M) . The following assertions are equivalent: (i) $\{x_t: t \in T\}$ is a system of mutually compatible observables.

(ii) $\{x_t: t \in T\}$ has a joint σ -observable.

(iii) There is a measure space (Ω, \mathcal{G}) , an \mathcal{G} -measurable function $g_t: \Omega \rightarrow R_1$, and an \mathcal{G} - σ -observable h of (X, M) such that $h(g_t^{-1}(E)) = x_t(E)$ for all $t \in T$ and $E \in B(R_1)$.

If, moreover, (X, M) is separable in the sense that any Boolean σ -algebra of (X, M) has a countable generator, or that T is coun-

table, then (i) is equivalent to (iv) There exists a \mathcal{G} -observable x and measurable functions $f_t: R_1 \rightarrow R_1$ such that, for all $t \in T$, $x_t = f_t \circ x$.

Proof. It follows from Theorem 2.1 and Theorem 6.9 of [9]. Q.E.D.

Theorem 2.3. Let x_1, \dots, x_n be mutually compatible \mathcal{G} -observables of an F quantum space (X, M) with a joint \mathcal{G} -observable x . If g is any real-valued Borel function on R_n , then $g \circ (x_1, \dots, x_n) : E \mapsto x(g^{-1}(E))$, $E \in B(R_1)$, is a \mathcal{G} -observable of (X, M) . If g_1, \dots, g_k are real-valued Borel functions on R_n and $y_i = g_i \circ (x_1, \dots, x_n)$, then y_1, \dots, y_k are mutually compatible \mathcal{G} -observables of (X, M) , and for any real-valued Borel function h on R_k $h \circ (y_1, \dots, y_k) = (h(g_1, \dots, g_k)) \circ (x_1, \dots, x_n)$, where $h(g_1, \dots, g_k)$ is the function $t = (t_1, \dots, t_n) \mapsto h(g_1(t), \dots, g_k(t))$.

Proof. It is straightforward and therefore is omitted. Q.E.D.

Theorems 2.2 and 2.3 are of a great importance for building so-called functional calculus for compatible \mathcal{G} -observables. Therefore, for compatible \mathcal{G} -observables x and y of (X, M) we may define $x+y$, $x \cdot y$, etc., if we put, for example, $x+y = (f+g) \circ z$, where $x = f \circ z$, $y = g \circ z$, according to Theorem 2.2, etc.

In the rest of this section we concentrate on the problem of existence of a joint distribution of a given system of mutually compatible \mathcal{G} -observables. For the quantum logic approach to quantum mechanics it is of great importance; it is known [2] that there are cases when it fails. By an F -state on F -quantum space (X, M) we understand a mapping $m: M \rightarrow [0, 1]$ such that $m(f \vee (1-f)) = 1$ for any $f \in M$; (ii) $m(\bigvee_i f_i) = \sum_i m(f_i)$ whenever $f_i \leq 1-f_j$ for $i \neq j$. In the terminology of Piasecki [4] and F -state is a P -measure.

Theorem 2.4. Let $\{x_t: t \in T\}$ be a system of \mathcal{G} -observables of an F -quantum space (X, M) . The pairwise compatibility of $\{x_t: t \in T\}$ implies that there exists a unique probability measure μ (called a joint distribution of $\{x_t: t \in T\}$) such that μ is called a joint distribution of $\{x_t: t \in T\}$ such that

$\mu(\prod_{t \in \mathcal{A}} A_t) = m(\bigwedge_{t \in \mathcal{A}} x_t(A_t))$ for any $A_t \in B(R_1)$, $t \in \mathcal{A}$, and any finite subset $\mathcal{A} \subset T$ in any F-state m .

Proof. According to Theorem 2.1, there is a joint \mathcal{G} -observable x of $\{x_t: t \in T\}$. Let us put $\mu(A) = m(x(A))$, $A \in B(T)$. We assert that $\mu = \mu^m$ is a probability measure in question. Q.E.D.

The authors hope to study also the problem of \mathcal{G} -joint distribution for noncompatible observables.

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